

Lieske and Null determined λ using isotropic coordinates and $\pm 0.08 < \lambda < 1.26 \pm 0.11$; based on the observations and using the Painlevé's $\lambda = 0.93 \pm 0.06$ (Sitarski 1983).

In the present work we used the Icarus and determined λ along with computing 1424 observational equations by the method. We obtained

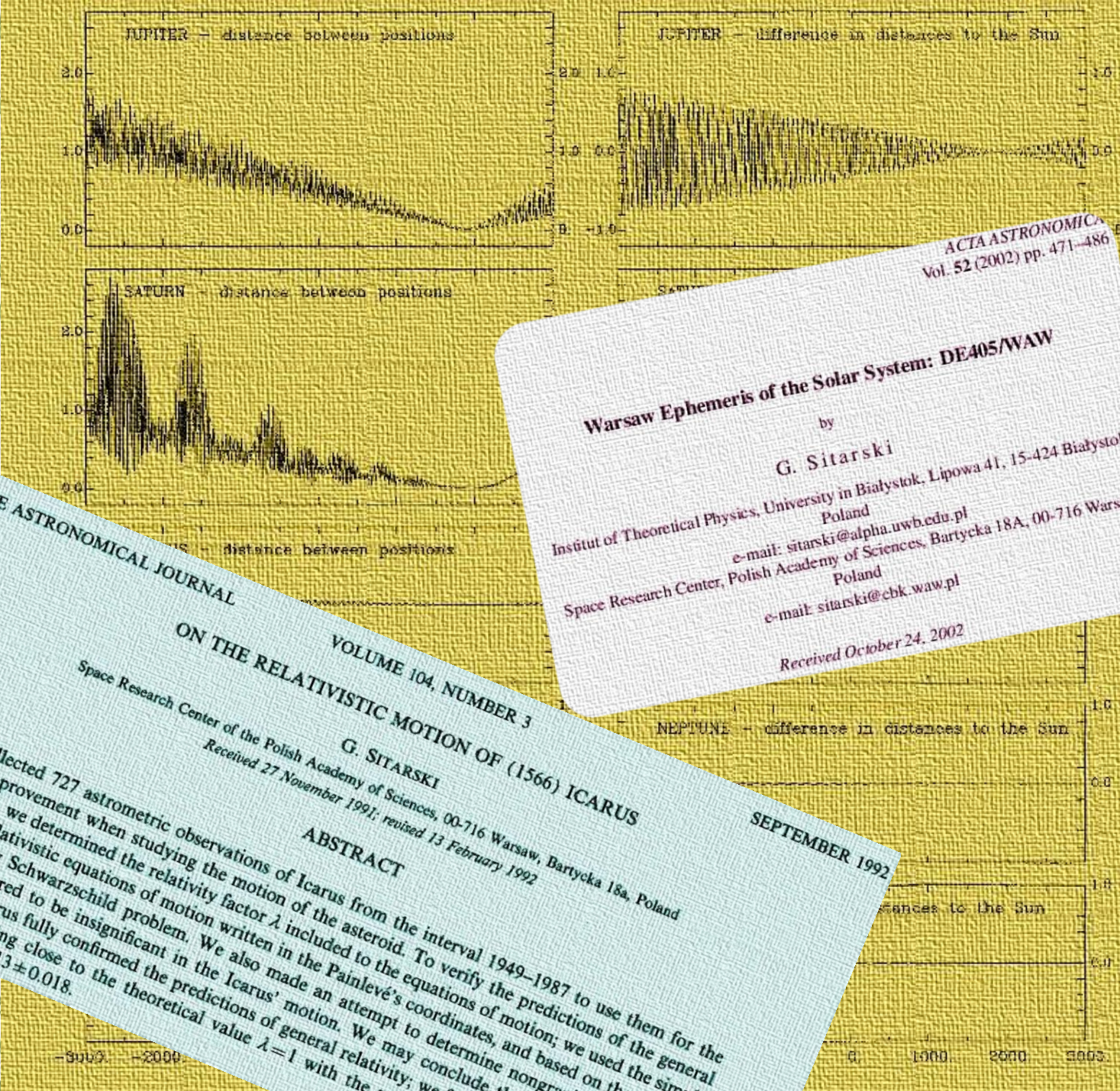
$$\lambda = 0.985 \pm 0.033.$$

We also made an attempt to take into account a hypothetical "nongravitational" effect in the form of a secular change $da/dt = \dot{a}$ of the semi-major axis. Hence, together with orbital elements a, e, i , and we got

$$\lambda = 1.013 \pm 0.018,$$

$$\dot{a} = (-0.206 \pm 0.109) \times 10^{-4} \text{ AU/yr}$$

We can see that in both cases (either with or without the nongravitational effect) the determined value of λ is very close to $\lambda = 1$, and a difference $|1 - \lambda|$ is comparable to the observational error.



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Warsaw Ephemeris of the Solar System: DE405/WAW
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ON THE RELATIVISTIC MOTION OF (1566) ICARUS
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ABSTRACT
We collected 727 astrometric observations of Icarus from the interval 1949–1987 to use them for the orbit improvement when studying the motion of the asteroid. To verify the predictions of the general relativity, we determined the relativity factor λ included to the equations of motion, and based on the simple form of relativistic equations of motion written in the Painlevé's coordinates, we used the simple form of one-body Schwarzschild equations of motion. We also made an attempt to determine nongravitational effects which appeared to be insignificant in the Icarus' motion. We may conclude that our analysis of the motion of Icarus fully confirmed the predictions of general relativity; we found a value of the relativity factor λ as being close to the theoretical value $\lambda = 1$ with the probability of 99.8%, namely $\lambda = 0.985 \pm 0.033 < \lambda < 1.013 \pm 0.018$.

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Generating of "Clones" of an Impact Orbit for the Earth-As Collision

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As like as in the above case, now let us define the functions

$$v = a(r/q)^{-\mu}, \quad w = (r/q)^{\nu}, \quad \psi = (1+w)^{-\kappa},$$

complete a set of the differential equations

$$r\dot{v} = -\mu v \dot{r}, \quad r\dot{w} = \nu w \dot{r}, \quad (1+w)\dot{\psi} = -\kappa \psi \dot{w},$$

and also write $g(r) = v\psi$. Thus we can again derive the recurrent relations to compute the values of coefficients g_n for the function $g(r)$:

$$\begin{aligned} v_0 &= a(r_0/q)^{-\mu}, & v_1 &= -\mu v_0 r_1/r_0, \\ w_0 &= (r_0/q)^{\nu}, & w_1 &= \nu w_0 r_1/r_0, \\ \psi_0 &= (1+w_0)^{-\kappa}, & \psi_1 &= -\kappa \psi_0 w_1/w_0, \\ g_0 &= v_0 \psi_0, \end{aligned}$$

and for $n = 1, \dots, N$

$$\begin{aligned} (n+1)r_0 v_{n+1} &= -\mu(n+1)v_0 r_{n+1} - \sum_{k=0}^{n-1} (k+1)(\mu r_k v_k), \\ (n+1)r_0 w_{n+1} &= \nu(n+1)w_0 r_{n+1} + \sum_{k=0}^{n-1} (k+1)(\nu r_k w_k), \\ (n+1)(1+w_0)\psi_{n+1} &= -\kappa(n+1)\psi_0 w_{n+1} - \sum_{k=0}^{n-1} (k+1)(\kappa w_k \psi_k), \\ g_n &= \sum_{k=0}^n v_k \psi_{n-k}. \end{aligned}$$

We can create six normal equations and then six elimination equations which can be schematically written as below:

normal equations						elimination equations							
Δx	Δy	Δz	$\Delta \dot{x}$	$\Delta \dot{y}$	$\Delta \dot{z}$	Δx	Δy	Δz	$\Delta \dot{x}$	$\Delta \dot{y}$	$\Delta \dot{z}$		
$[aa]$	$[ba]$	$[ca]$	$[da]$	$[ea]$	$[fa]$	$[la]$	$r_{a,1}$	$r_{b,1}$	$r_{c,1}$	$r_{d,1}$	$r_{e,1}$	$r_{f,1}$	ρ_1
$[ab]$	$[bb]$	$[cb]$	$[db]$	$[eb]$	$[fb]$	$[lb]$	$r_{b,2}$	$r_{c,2}$	$r_{d,2}$	$r_{e,2}$	$r_{f,2}$		ρ_2
$[ac]$	$[bc]$	$[cc]$	$[dc]$	$[ec]$	$[fc]$	$[lc]$	$r_{c,3}$	$r_{d,3}$	$r_{e,3}$	$r_{f,3}$			ρ_3
$[ad]$	$[bd]$	$[cd]$	$[dd]$	$[ed]$	$[fd]$	$[ld]$	$r_{d,4}$	$r_{e,4}$	$r_{f,4}$				ρ_4
$[ae]$	$[be]$	$[ce]$	$[de]$	$[ee]$	$[fe]$	$[le]$	$r_{e,5}$	$r_{f,5}$					ρ_5
$[af]$	$[bf]$	$[cf]$	$[df]$	$[ef]$	$[ff]$	$[lf]$	$r_{f,6}$						ρ_6

what means that e.g.:

$$[aa]\Delta x + [ba]\Delta y + [ca]\Delta z + [da]\Delta \dot{x} + [ea]\Delta \dot{y} + [fa]\Delta \dot{z} = [la],$$

$$r_{a,1}\Delta x + r_{b,1}\Delta y + r_{c,1}\Delta z + r_{d,1}\Delta \dot{x} + r_{e,1}\Delta \dot{y} + r_{f,1}\Delta \dot{z} = \rho_1.$$

Values of $r_{a,1}, r_{b,1}, \dots, \rho_1$, and further $r_{b,2}, r_{c,2}, \dots, \rho_2$ and so on are found using earlier computed values of $[aa], [ba], \dots, [fa]$. Tables of normal and of elimination equations are given in the Appendix. We observe that the elimination equations are linearly independent.

Recurrent Power Series Integration of Equations of Comet's Motion Including the Nongravitational Terms in Marsden's Form

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ABSTRACT
Recurrent relations for computing the coefficients in time power expansion of $f(r)$ and $g(r)$ are derived what makes possible to integrate the equations of motion including the nongravitational forces in the form $A_j f(r)$ or $A_j g(r)$ are also given to compute the coefficients of the power series of $f(r)$ and $g(r)$ by the least squares method.